# **Gravity Tree Amplitudes at Infinity**

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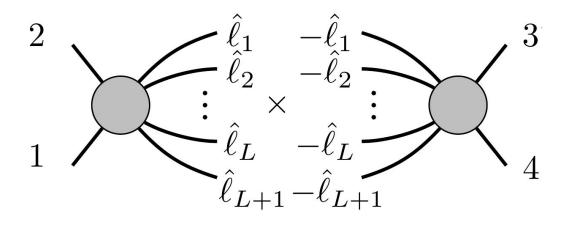
## Cut integrands, chiral shift

Take an L-loop cut 4-point amplitude and deform momenta by:

$$egin{aligned} \widetilde{\lambda}_{\ell_k} & \longrightarrow & \widehat{\widetilde{\lambda}}_{\ell_k} = \widetilde{\lambda}_{\ell_k} + z\,lpha_k\,\widetilde{\eta} & \sum_{k=1}^{L+1} lpha_k\,\lambda_{\ell_k} = 0 \end{aligned}$$

$$\sum_{k=1}^{L+1} lpha_k \, \lambda_{\ell_k} = 0$$

Integrand scaling is improved in d = 4



$$\sim rac{1}{z^5} \quad (\mathcal{N}=8)$$

(Edison, Hermann, Parra-Martinez, Trnka 2019)

## 'Poles' at infinity

- Study leading behavior of amplitudes at infinity
- Not usual  $p_i^2 = 0$  poles, limit is ambiguous
  - No true factorization
- Shift is a prescription to take the limit
- Different z scalings

## Leading 6pt term from chiral shift

• Example of a leading term at  $O(z^2)$ 

$$egin{aligned} \mathcal{M}_6(1^-2^+\hat{3}^-\hat{4}^-\hat{5}^+\hat{6}^+) &\sim z^2 rac{[2\eta]^8}{[12]^2[1\eta]^2[2\eta]^2} rac{\langle 34
angle^8}{\langle 35
angle\langle 46
angle\langle 3 ilde{\eta}]\langle 4 ilde{\eta}]\langle 5 ilde{\eta}]\langle 6 ilde{\eta}]} \ & imes \left(rac{[\chi_{36}\eta][\chi_{45}\eta]}{\langle 36
angle\langle 45
angle} + rac{[\chi_{34}\eta][\chi_{56}\eta]}{\langle 34
angle\langle 56
angle}
ight) + \mathcal{O}(z^2) \end{aligned}$$

Unshifted legs do not mix with shifted

#### Pole for *n*-point amplitude

Pattern continues for *n*-point amplitude:

$$egin{split} \mathcal{M}_n (1^-2^+\hat{3}^+\hat{4}^-5^+\dots\,\hat{n}^+) \sim & (-1)^n z^{n-4} s_{12} \mathcal{M}_3 (1^-2^+\eta^+) \ & imes \mathcal{M}_{n-1} (3^-4^-5^+\dots n^+\omega^+) + \mathcal{O}(z^{n-5}) \end{split}$$

Similar formulas exist for other helicity choices

#### Other patterns

More legs unshifted or different helicities have different forms:

$$\mathcal{M}_6(1^-2^-3^+\hat{4}^+\hat{5}^+\hat{6}^+)\sim -z^3rac{\langle 12
angle^8[1n][2n][3\eta]lpha_4lpha_5lpha_6}{\langle 14
angle\langle 15
angle\langle 16
angle\langle 24
angle\langle 25
angle\langle 26
angle\langle 34
angle\langle 35
angle\langle 36
angle}$$

Leading behavior would vanish due to helicities

### Behavior of YM Amplitudes

Near-identical 'factorization' to the GR case

$${\cal A}(1^-2^+\hat{3}^-\hat{4}^-\hat{5}^+\hat{6}^+) \sim -z^0 rac{[2\eta]^4}{[12][1\eta][2\eta]} rac{\langle 34 
angle^4}{\langle 3 ilde{\eta}] \langle 34 
angle \langle 45 
angle \langle 56 
angle \langle 6 ilde{\eta}]}$$

• Similar *n*-point formula

$${\cal A}(1^-2^+\hat{3}^-\hat{4}^-\hat{5}^+\cdots\hat{n}^+) \sim -z^0{\cal A}(1^-2^+\eta^+){\cal A}(3^-4^-5^+\cdots n^+\omega^+)$$

#### **Outlook and goals**

- Understand UV behavior on cut using tree amplitudes at infinity
- Catalogue patterns present at infinity
- Write *n*-point formulas for all possible helicities

Thank you for your time!

