

# MHV Amplitudes on Self-Dual Plane Waves

Tim Adamo  
University of Edinburgh

QCD Meets Gravity

12 December 2019

Work in progress with Lionel Mason & Atul Sharma  
(also work with E. Casali, A. Ilderton, S. Nekovar)

# Motivation

Many reasons to be interested in perturbative QFT in *strong* (non-trivial) background fields

- Practical (strong field QED/QCD, cosmology, GWs, holography, non-perturbative effects)
- Theoretical (probe robustness of structures in pQFT)

# Motivation

Many reasons to be interested in perturbative QFT in *strong* (non-trivial) background fields

- Practical (strong field QED/QCD, cosmology, GWs, holography, non-perturbative effects)
- Theoretical (probe robustness of structures in pQFT)

Challenges cut across both!

Example: tree-level gauge theory and gravity

- Flat background: full tree-level S-matrix
- Even simplest strong backgrounds: only 3- or (sometimes) 4-points

# Today

Is there any hope to make all-multiplicity statements on strong backgrounds?

# Today

Is there any hope to make all-multiplicity statements on strong backgrounds?

YES!

- Parke-Taylor-like formula for MHV gluon scattering on a self-dual plane wave background

$$\delta_{+,\perp}^3 \left( \sum_{i=1}^n k_i \right) \frac{\langle ij \rangle^4}{\langle 12 \rangle \langle 23 \rangle \cdots \langle n1 \rangle} \int_{-\infty}^{+\infty} dx^- e^{i\mathcal{F}_n(x^-)}$$

# Today

Is there any hope to make all-multiplicity statements on strong backgrounds?

YES!

- Parke-Taylor-like formula for MHV gluon scattering on a self-dual plane wave background

$$\delta_{+,\perp}^3 \left( \sum_{i=1}^n k_i \right) \frac{\langle ij \rangle^4}{\langle 12 \rangle \langle 23 \rangle \cdots \langle n1 \rangle} \int_{-\infty}^{+\infty} dx^- e^{i\mathcal{F}_n(x^-)}$$

- Full tree-level S-matrix conjecture for Yang-Mills on such backgrounds

# Plane waves

Solution to vacuum equations (in  $d$  dim.) with:

- covariantly constant null symmetry  $n$ ,
- $(2d - 4)$  additional symmetries,
- commuting to form Heisenberg algebra w/ center  $n$

# Plane waves

Solution to vacuum equations (in  $d$  dim.) with:

- covariantly constant null symmetry  $n$ ,
- $(2d - 4)$  additional symmetries,
- commuting to form Heisenberg algebra w/ center  $n$

For Yang-Mills theory, PWs valued in Cartan of gauge group

[Trautman, Basler-Hadicke, TA-Casali-Mason-Nekovar]

$$ds^2 = 2dx^+ dx^- - (dx^\perp)^2, \quad A = x^\perp \dot{a}_\perp(x^-) dx^-$$



# Plane waves

Solution to vacuum equations (in  $d$  dim.) with:

- covariantly constant null symmetry  $n$ ,
- $(2d - 4)$  additional symmetries,
- commuting to form Heisenberg algebra w/ center  $n$

For Yang-Mills theory, PWs valued in Cartan of gauge group

[Trautman, Basler-Hadicke, TA-Casali-Mason-Nekovar]

$$ds^2 = 2dx^+ dx^- - (dx^\perp)^2, \quad A = x^\perp \dot{a}_\perp(x^-) dx^-$$

$\dot{a}_\perp(x^-)$  compactly supported  $\leftrightarrow$  well-defined S-matrix [Schwinger,

TA-Casali-Mason-Nekovar]

## Self-dual plane waves

In  $d = 4$ , complexify  $\mathbb{R}^{1,3}$  to  $\mathbb{C}^4$ :  $ds^2 = 2(dx^+ dx^- - dz d\bar{z})$ .  
Require PW & self-duality:

$$*F = iF$$

## Self-dual plane waves

In  $d = 4$ , complexify  $\mathbb{R}^{1,3}$  to  $\mathbb{C}^4$ :  $ds^2 = 2(dx^+ dx^- - dz d\tilde{z})$ .  
Require PW & self-duality:

$$*F = iF$$

Propagation direction of wave:  $n = \frac{\partial}{\partial x^+}$

$$\text{Since } n^2 = 0, \quad n^{\alpha\dot{\alpha}} = \iota^\alpha \tilde{\iota}^{\dot{\alpha}}, \quad \iota^\alpha = \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \tilde{\iota}^{\dot{\alpha}}$$

Result:

$$A = \tilde{z} \dot{f}(x^-) dx^- = \tilde{z} \dot{f}(x^-) \iota_\alpha \tilde{\iota}^{\dot{\alpha}} dx^{\alpha\dot{\alpha}}$$
$$F = \dot{f}(x^-) d\tilde{z} \wedge dx^- = \dot{f} \tilde{\iota}_{\dot{\alpha}} \tilde{\iota}_{\dot{\beta}} dx_\alpha^{\dot{\alpha}} \wedge dx^{\alpha\dot{\beta}}$$

# SDPW kinematics

The spinor-helicity formalism works on *all* plane waves [TA-Ilderton]

SDPW have *chiral* on-shell kinematics

## SDPW kinematics

The spinor-helicity formalism works on *all* plane waves [TA-Ilderton]

SDPW have *chiral* on-shell kinematics

Gluon with incoming momentum  $k_{\alpha\dot{\alpha}} = \lambda_{\alpha}\tilde{\lambda}_{\dot{\alpha}}$ :  $T^a \mathcal{E}_{\alpha\dot{\alpha}}^{\pm}(x^-) e^{i\phi_k}$

$$\phi_k = k \cdot x + e\tilde{z} f(x^-) + \frac{k}{k_+} \int^{x^-} dt e f(t)$$

On-shell kinematics:

$$K_{\alpha\dot{\alpha}}(x^-) = \lambda_{\alpha} \tilde{\Lambda}_{\dot{\alpha}}, \quad \tilde{\Lambda}_{\dot{\alpha}} := \tilde{\lambda}_{\dot{\alpha}} + \frac{e}{\sqrt{k_+}} \tilde{\iota}_{\dot{\alpha}} f(x^-)$$

$$\mathcal{E}_{\alpha\dot{\alpha}}^- = \frac{\lambda_{\alpha} \tilde{\iota}_{\dot{\alpha}}}{[\tilde{\iota} \tilde{\lambda}]}, \quad \mathcal{E}_{\alpha\dot{\alpha}}^+ = \frac{\iota_{\alpha} \tilde{\Lambda}_{\dot{\alpha}}}{\langle \iota \lambda \rangle}$$

# Twistor theory

Twistor space:  $Z^A = (\mu^{\dot{\alpha}}, \lambda_{\alpha})$  homog. coords. on  $\mathbb{CP}^3$

$$\mathbb{PT} = \mathbb{CP}^3 \setminus \{\lambda_{\alpha} = 0\}$$

$x \in \mathbb{C}^4$  given by  $X \cong \mathbb{CP}^1 \subset \mathbb{PT}$  via  $\mu^{\dot{\alpha}} = x^{\alpha\dot{\alpha}} \lambda_{\alpha}$

# Twistor theory

Twistor space:  $Z^A = (\mu^{\dot{\alpha}}, \lambda_{\alpha})$  homog. coords. on  $\mathbb{CP}^3$

$$\mathbb{PT} = \mathbb{CP}^3 \setminus \{\lambda_{\alpha} = 0\}$$

$x \in \mathbb{C}^4$  given by  $X \cong \mathbb{CP}^1 \subset \mathbb{PT}$  via  $\mu^{\dot{\alpha}} = x^{\alpha\dot{\alpha}} \lambda_{\alpha}$

Familiar applications from flat background:

- Massless free fields  $\leftrightarrow$  cohomology on  $\mathbb{PT}$  [Penrose, Sparling, Eastwood-Penrose-Wells]
- Representation for on-shell scattering kinematics [Hodges]
- Full tree-level S-matrix of  $\mathcal{N} = 4$  SYM [Witten, Berkovits, Roiban-Spradlin-Volovich]
- Full tree-level S-matrix of  $\mathcal{N} = 8$  SUGRA [Cachazo-Skinner]

What's this got to do with perturbation theory on SDPWs?



What's this got to do with perturbation theory on SDPWs?

## Theorem [Ward, 1977]

There is a 1:1 correspondence between:

- SD  $SU(N)$  Yang-Mills fields on  $\mathbb{C}^4$ , and
- rank  $N$  holomorphic vector bundles  $E \rightarrow \mathbb{P}^1$  trivial on every  $X \subset \mathbb{P}^1$  (+ technical conditions)

What's this got to do with perturbation theory on SDPWs?

## Theorem [Ward, 1977]

There is a 1:1 correspondence between:

- SD  $SU(N)$  Yang-Mills fields on  $\mathbb{C}^4$ , and
- rank  $N$  holomorphic vector bundles  $E \rightarrow \mathbb{P}^1$  trivial on every  $X \subset \mathbb{P}^1$  (+ technical conditions)

**Upshot:** twistor theory *trivializes* the SD sector

## SDPWs in Twistor Space

Can construct  $E \rightarrow \mathbb{P}\mathbb{T}$  explicitly; holomorphicity encoded by partial connection on  $E$ :

$$\bar{D} = \bar{\partial} + A, \quad A = \int_{\mathbb{C}^*} \frac{ds}{s} \bar{\delta}^2(\iota - s\lambda) \int^{s[\tilde{\iota}\mu]} dt f(t)$$

Easy to show that  $\bar{D}^2 = 0$

## SDPWs in Twistor Space

Can construct  $E \rightarrow \mathbb{P}\mathbb{T}$  explicitly; holomorphicity encoded by partial connection on  $E$ :

$$\bar{D} = \bar{\partial} + A, \quad A = \int_{\mathbb{C}^*} \frac{ds}{s} \bar{\delta}^2(\iota - s\lambda) \int^{s[\tilde{\iota}\mu]} dt f(t)$$

Easy to show that  $\bar{D}^2 = 0$

**Penrose transform:** gluons encoded by  $E$ -twisted cohomology on  $\mathbb{P}\mathbb{T}$

$$\text{- helicity} \leftrightarrow H_{\bar{D}}^{0,1}(\mathbb{P}\mathbb{T}, \mathcal{O}(-4) \otimes E)$$

$$\text{+ helicity} \leftrightarrow H_{\bar{D}}^{0,1}(\mathbb{P}\mathbb{T}, \mathcal{O} \otimes E)$$

So what?

So what?

Can give perturbative formulation of SDPW background field  
Yang-Mills in  $\mathbb{P}^1$  [Mason, Boels, TA-Mason-Sharma]

Properties:

- Perturbative around SD sector
- Generating functional for MHV interactions localized on lines  $X \subset \mathbb{P}^1$
- Explicit expansion using holomorphic trivialization of  $E|_X$

# MHV amplitude

Evaluating this expansion gives:

$$\delta_{+,\perp}^3 \left( \sum_{i=1}^n k_i \right) \frac{\langle ij \rangle^4}{\langle 12 \rangle \langle 23 \rangle \cdots \langle n1 \rangle} \int_{-\infty}^{+\infty} dx^- e^{i\mathcal{F}_n(x^-)}$$

for *Volkov exponent*

$$\mathbb{K}^{\alpha\dot{\alpha}}(x^-) := \sum_{i=1}^{n-1} K_i^{\alpha\dot{\alpha}}(x^-),$$

$$\mathcal{F}_n(x^-) := \frac{1}{\mathbb{K}_+} \int^{x^-} dt \mathbb{K}^2(t)$$

Twistor action proves formula is correct, but...



Twistor action proves formula is correct, but...

**Red flag:** only *one* residual lightfront integral!

Expect  $n - 2$  for  $n$ -point tree amplitude

Twistor action proves formula is correct, but...

**Red flag:** only *one* residual lightfront integral!

Expect  $n - 2$  for  $n$ -point tree amplitude

**Resolution:** field redefinition recasts Yang-Mills action such that all MHV vertices have *single* lightfront integral [Mansfield]

Twistor action proves formula is correct, but...

**Red flag:** only *one* residual lightfront integral!

Expect  $n - 2$  for  $n$ -point tree amplitude

**Resolution:** field redefinition recasts Yang-Mills action such that all MHV vertices have *single* lightfront integral [Mansfield]

Other sanity checks & features:

- Explicit checks at 3- and 4-points
- Perturbative limit ( $\text{MHV}_n + \text{background} \rightarrow \text{MHV}_{n+1}$ )
- Flat background limit
- Generalization to  $\mathcal{N} = 4$  SYM

## Full tree-level S-matrix?

Easy guess for  $N^k$ MHV, based on holomorphic maps

$$Z : \mathbb{CP}^1 \rightarrow \mathbb{PT}$$

$$\int \frac{\prod_{a=0}^{k+1} d^{4|4} U_a}{\text{vol GL}(2, \mathbb{C})} \text{tr} \left( \prod_{i=1}^n \frac{\gamma_i^{-1} \mathcal{A}_i \gamma_i d\sigma_i}{\sigma_i - \sigma_{i+1}} \right)$$

where:

- $Z(\sigma) = \sum_{a=0}^{k+1} U_a \sigma^a$  is a degree  $k + 1$  holomorphic map
- $\{\sigma_i\} \subset \mathbb{CP}^1$  punctures on  $\mathbb{CP}^1$
- $\mathcal{A} \in H_D^{0,1}(\mathbb{PT}, \mathcal{O} \otimes E)$  twistor wavefunctions
- $\gamma$  holomorphic frame trivializing  $E$  over image of  $Z$

## Full tree-level S-matrix?

Easy guess for  $N^k$ MHV, based on holomorphic maps

$$Z : \mathbb{CP}^1 \rightarrow \mathbb{PT}$$

$$\int \frac{\prod_{a=0}^{k+1} d^{4|4} U_a}{\text{vol GL}(2, \mathbb{C})} \text{tr} \left( \prod_{i=1}^n \frac{\gamma_i^{-1} \mathcal{A}_i \gamma_i d\sigma_i}{\sigma_i - \sigma_{i+1}} \right)$$

where:

- $Z(\sigma) = \sum_{a=0}^{k+1} U_a \sigma^a$  is a degree  $k + 1$  holomorphic map
- $\{\sigma_i\} \subset \mathbb{CP}^1$  punctures on  $\mathbb{CP}^1$
- $\mathcal{A} \in H_D^{0,1}(\mathbb{PT}, \mathcal{O} \otimes E)$  twistor wavefunctions
- $\gamma$  holomorphic frame trivializing  $E$  over image of  $Z$

Currently just a conjecture...

# Summary

**Upshot:** it *is* possible to make all-multiplicity statements in strong backgrounds!

Many exciting things to do:

- Prove/correct  $N^k$ MHV conjecture
- Gravitational SDPWs
- Double copy for full tree-level SDPW S-matrix
- Generalize to generic PW backgrounds